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# STUDIES ON ALTERNATING CURRENT ELECTROLYTIC CELLS

Final Report to the  
OFFICE OF NAVAL RESEARCH  
June 15, 1954

Research Contract  
Nonr-604(00)

Project Number  
NR-051-288

By  
A. Edward Remick, Herbert W. McCormick and William K. Sneed  
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Detroit, Michigan

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## STUDIES ON ALTERNATING CURRENT

### ELECTROLYTIC CELLS

#### PART I

#### The Effects of Concentration on Polarization Capacity and Polarization Resistance in Ferrocyanide-Ferricyanide Cells with Platinum Electrodes

By: A. Edward Remick and Herbert W. McCormick

The results of this research have already been reported in our Technical Reports No. 1 and No. 2. There is no need, therefore, to do more here than outline the objectives, the method used and the conclusions. Literature references will be given by using the Roman numerals I and II to refer respectively to our Technical Reports No. 1 and No. 2 and Arabic numerals will be used to indicate the reference numbers used in our previous bibliographies.

#### Objectives

Our ultimate objective is three fold. (1) We wish to analyse the behavior of an electrode in an alternating current electrolytic cell into its component properties. The result of such an analysis can be represented graphically by an equivalent circuit consisting of both conventional and unconventional circuit elements. (2) We wish to establish the fundamental nature of the unconventional circuit elements. (3) We wish to apply this fundamental knowledge to the elucidation of the mechanism of electrochemical redox reactions.

An excellent start toward our goal has been made by Rozental and Ershler (II-2), Randles (II-3,7), Grahame (II-4)

and Gerischer (II-5). Their mathematical theories and equivalent circuits are closely similar but Grahame's theory is the most complete and we have therefore concentrated our efforts on collecting experimental data with which to check his theory. Specifically we have studied the influence of frequency, depolarizer concentration and ionic strength on the equivalent series capacitance ( $C_s$ ) and resistance ( $R_s$ ) of an alternating current electrolytic cell and on the equivalent series capacitance ( $C_s^*$ ) and resistance ( $R_s^*$ ) of the faradaic branch of the equivalent circuit of said cell. These measurements were made using very small current densities. A little work was also done at high current densities for the purpose of furthering the studies of Shaw and Remick (I-1) which give us valuable qualitative "tools" for the investigation of electrode mechanism.

#### Experimental Method

The electrolytic cell was fitted with platinum electrodes arranged so that the interelectrode distance could be varied and measured. A piece of glass tubing was slipped over the two electrodes with the upper end projecting above the surface of the liquid; this was done to insure a uniform cross-sectional area of the conducting path. The electrolyte always contained equimolar concentrations of potassium ferrocyanide and ferricyanide and usually contained also potassium sulfate or sodium benzene sulfonate as supporting electrolytes; it was always freed of dissolved oxygen by a stream of pure

nitrogen. Sinusoidal alternating current was passed through the cell. Measurements of resistance and capacitance were made on an impedance bridge using an oscillograph for a detecting instrument. Measurements were carried out at  $30^{\circ} \pm 0.1^{\circ}$ .

### Conclusions

- (1) Reproducible results can be obtained with platinum electrodes if they are cleaned by a series of alternate cathodic and anodic electrolyses in an alkaline cyanide solution.
- (2) Hysteresis is observed at sufficiently high current densities. It's probable cause is the collection of electrolytic gases on the surfaces of the electrodes.
- (3) When  $C_s$  and  $R_s$  are plotted against current density, a "capacitive depression point" and a "resistive elevation point" are observed at the same current density. An explanation readily emerges from Grahame's theory.
- (4) There is a frequency-independent component of the cell resistance in addition to the electrolytic resistance ( $R_T$ ). We tentatively call it "the electrode layer resistance" and symbolize it as  $R_L$ . We may therefore write
$$R_s = R_T + R_L + R_{\Delta} \quad (1)$$
where  $R_L$  is the frequency dependent component of  $R_s$ .
- (5)  $R_L$  is always positive. It increases as the depolarizer concentration increases and decreases as the concentration of the supporting electrolyte increases, other things

being equal.

- (6) Our improved method of cleaning electrodes makes a constant contribution to  $R_L$  whereas four other cleaning methods gave four values of  $R_L$ , other things being equal. In the latter case, a rough parallelism was observed between  $R_L$  and  $C_s$ .
- (7)  $C_s$  increases with the ionic strength slightly and with the depolarizer concentration markedly.
- (8) Experiments made with potassium sulfate and sodium benzenesulfonate as supporting electrolytes showed that the quantity  $(R_L + R_A)$  is accurately the same at a given value of the ionic strength independent of the kind of supporting electrolyte used, other things being equal. This indicates that the effect of the supporting electrolyte on  $(R_A + R_L)$  is an ionic strength effect, not a mere concentration effect.  
  
(Note: The values of  $R_A$  listed in Tables I-IV of our Technical Report No. 1 should be interpreted as  $R_{LA} + R_L$  in terms of our present definition of  $R_A$ .)
- (9) The quantity  $(R_A + R_L)$  decreases markedly as the ionic strength increases and is essentially independent of depolarizer concentration.
- (10)  $C_s$  is a linear function of the depolarizer concentration in the presence of sufficient supporting electrolyte. The slopes of these curves are positive and are greater at lower frequencies, the half-wave area of the alter-

nating current being constant. The same conclusion may be drawn about  $C_S^*$ .

- (11) Grahame's theory leads to the equations

$$R_S^* = \eta / \sqrt{\omega} \quad (2)$$

$$C_S^* = 1/\eta \sqrt{\omega} \quad (3)$$

Eq.(2) involves the assumption that Grahame's  $\theta$  is zero. This is equivalent to the assumption that the electron transfer reactions are rapid compared to the period of the alternating current. Plots of the experimental values of  $R_S^*$  and  $C_S^*$  against  $1/\sqrt{\omega}$  were found to be linear passing thru the origin, as demanded by Eq's. (1) and (2). Moreover, values of  $\eta$  calculated from the slopes of the  $C_S^*$  and  $R_S^*$  curves agreed fairly well and the equivalent requirement that  $C_S^* R_S^* \omega = 1$  was accurately satisfied at frequencies up to 1000 c.p.s. if a sufficiently large excess of supporting electrolyte is used. The dispersion studies here involved were made at constant half-wave area of current, i.e., they were "Q-dispersion" studies.

- (12) From the standpoint of the effect of depolarizer concentration (C) on  $C_S^*$  and  $R_S^*$ , Grahame's theory leads to the equations:

$$C_S^* = kC \quad (4)$$

$$R_S^* = k'/C \quad (5)$$

if we make the assumptions that: (a) equilibrium in the electron transfer reaction is continuously established



between the electrode and the adjacent layer of electrolyte, and (b) using time-average values, the depolarizer concentrations at the electrode surface are the same as those in the bulk of the solution. Eq's. (4) and (5) demand a linear dependence between  $C_s^*$  and  $1/R_s^*$ , on the one hand, with  $C$  on the other hand. The predicted relations are found experimentally if a sufficient excess of supporting electrolyte is used.

The relationship between the constants of Eq's. (4) and (5) is:

$$k = 1/k'w$$

The experimental data is also in accord with this theoretical requirement.

- (13) The concordance between theory and experiment cited in conclusion #12 was achieved only by taking  $R_L$  into account and thus furnishes additional evidence in favor of the reality of  $R_L$ .
- (14) Consideration was given to the possibility that the resistance we have called  $R_L$  might more correctly be ascribed to a slow electron discharge step or, in other words, that  $R_L = 0$  and  $\theta$  is finite. It was pointed out that the oscillographic studies made by Silverman and Remick (II-10) on the ferricyanide-ferrocyanide system furnished evidence that  $\theta$  is essentially zero. Furthermore, our method of evaluating  $R_L$  involved extrapolation of  $R_s$  to infinite frequency. In terms of Grahame's

theory the impedance of the faradaic branch of Grahame's equivalent circuit, including  $\theta$ , is infinite when the frequency is infinite. It follows that the measured resistance which we have symbolized as  $R_L$  cannot be in the faradaic branch of the circuit and hence cannot be either wholly or partially accounted for as a resistance associated with the activation process, i.e., as  $\theta$ .

- (15) The accord achieved between experiment and Grahame's theory, using Grahame's equivalent circuit modified by the introduction of  $R_L$ , makes clear the fundamental nature of each of the circuit elements except  $R_L$ . Their nature is adequately described by Grahame's mathematical theory.
- (16) An increase in ionic strength at constant concentration of depolarizer and constant frequency results in an increase of  $C_S^*$  while  $R_S^*$  usually rises slightly to a maximum and then falls comparatively rapidly. The maximum in the  $R_S^*$  curve may not be real.

There are apparent three ways in which the ionic strength might affect the values of  $C_S^*$  and  $R_S^*$ :

- (a) Up to a point, an increase in ionic strength would diminish the effect of electrical migration. Let us call this the migration effect. (b) The coulombic interactions between the electrode and the depolarizer ions would be diminished by an increase in ionic strength in a manner reminiscent of the Brønsted-Christiansen-Scatchard equation

used in chemical kinetics. Let us call this the kinetic effect. (c) The activity coefficients of the depolarizer ions would be affected by interionic attractive forces and thus changes in ionic strength would alter the Nernst potentials which are associated with Grahame's and hence with  $C_S^*$  and  $R_S^*$  by Eq's. (2) and (3). Let us call this the thermodynamic effect.

We carried out a mathematical development based on the Debye-Hückel limiting law and Grahame's theory, which showed that the thermodynamic effect would operate in the opposite direction to that experimentally observed. The kinetic effect cannot be operative because  $\theta$  is zero. Thus if any one of these three effects is the major cause of the observed influence of the ionic strength on the values of  $C_S^*$  and  $R_S^*$ , it must be the migration effect. It is possible, at least in principle, to expand Grahame's mathematical treatment to include the migration effect. This we have not yet attempted. It must be done before we can know whether or not the influence of the ionic strength observed experimentally can be explained on this basis.

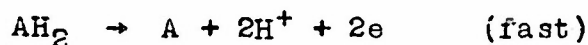
## PART II

### The Determination of Thermodynamic Reduction Potentials by an Alternating Current Technique

By: William K. Snead and A. Edward Remick

We have submitted no Technical Report on this research work because no phase of it has been successfully concluded. This report will deal, therefore, only with the guiding principles we have employed, with the unsuccessful laboratory methods we have tried and with the method we are currently developing.

The objective of this phase of our research was a method for measuring the thermodynamic reduction potentials of reversible electron-transfer steps occurring in redox reactions whose overall behavior would be described as irreversible. An example would be a redox reaction having a Conant mechanism:



In this case the problem would be to measure the reduction potential of the first step which involves the unstable intermediate free radical A.

The method we originally used was essentially the method of Silverman and Remick (I-17). We elected to study first reversible redox systems whose potentials could be measured potentiometrically, in this way hoping to establish the validity of the method. This done, we reasoned that the method could then be applied to irreversible oxidation or

reduction reactions to give the potential of the reversible electron transfer step, if any.

The electrical circuit we employed was the one used by Silvermand and Remick. A sinusoidal alternating current was passed through a cell containing in aqueous solution one component of a reversible redox system. One electrode of this cell was of smooth platinum sealed in glass so that no edges were exposed and had an exposed surface area of a few tenths of a square centimeter. The instantaneous potential of this electrode (under the influence of the exciting current) was impressed on the vertical plates of a DC oscilloscope, a standard electrode (saturated calomel electrode) and salt bridge completing the circuit to the electrolyte. Of course, none of the impressed AC current flowed in this oscilloscope-standard electrode-salt bridge arm of the circuit. The other electrode of the working cell was in most cases somewhat larger than the working electrode (about 2-3 cm<sup>2</sup>) and sometimes platinized.

The instantaneous potential at either of the two points of an oscillogram where the current is zero we have designated as  $e_0$ . Many years ago Reichenstein observed that as the r.m.s. current through a cell in which water was being electrolysed was increased, both the cathodic and anodic values of  $e_0$  (measured between the two working electrodes) approached limiting values. The difference between the two limiting values he considered as being equal to the reversible potential difference between a hydrogen and an oxygen

electrode. Silverman and Remick showed, however, that this potential difference included the corresponding overvoltages. We presumed that when reversible systems were used, and  $e_0$  corresponded to the potential difference between the working electrode and a reference half cell, the maximum value of the anodic  $e_0$  (in case the electrolyte was the reduced form of the redox system) should be the standard potential of the redox system against the given reference electrode, much as Reichenstein had originally imagined. This presumption was supported by the observation of Silverman and Remick that as the current was increased the oscillogram approached a condition wherein the part of the curve passing through the  $e_0$  value in question was linear for a considerable distance on each side of the  $e_0$  point. They advanced qualitative evidence to show that this linear region appeared when the current excursions became large enough to carry the system into the poised region. Thus the maximum value of  $e_0$  should occur when the redox ratio at the electrode surface is unity. Since the IR drop is zero at this point,  $e_0$  should be equal to the standard potential.

Silverman and Remick were unable to test their theory quantitatively because they did not have a D.C. oscillograph. We therefore undertook the test using a Dumont type 304-H oscillograph and a Dumont oscillographic record camera. As depolarizers we used potassium ferrocyanide, hydroquinone, stannous chloride and ferrous sulfate respectively in four different groups of experiments. In order to locate the

maximum value of  $e_0$  as accurately as possible we plotted  $e_0$  against the r.m.s. current. In all cases the values of  $e_0$  tended to level off but they never reached a maximum value or usually a point of inflection. They leveled off at values of the voltage close to the expected standard potential but obviously it was impossible to locate the point precisely. Some inflections were observed, however, but these were never reproducible nor retraceable because they occurred in the range of current density where hysteresis was operative.

The above work was done prior to the advent of Grahame's theory of the faradaic admittance (I-3). When we became acquainted with that theory it became evident to us that when the cell current ( $i_T$ ) is zero, it by no means follows that the faradaic current ( $i_F$ ) is also zero; hence our  $e_0$  undoubtedly contained an  $iR$  drop. We thereupon set to work trying to devise a method, both experimental and mathematical, by which we could locate the point at which the redox ratio is unity and correct the corresponding potential for the  $iR$  drop in the faradaic part of the circuit. We devised an experimental means of automatically subtracting the  $i_T R_T$  drop (where  $R_T$  is the electrolytic resistance) before the voltage was recorded on the screen of the oscillograph. We devised a number of different circuits involving square waves, triangular waves or sine waves to be used with the oscillograph and worked out their mathematical properties only to find that none of them gave unequivocal answers to our problem.

Incorporation of the effects of the electrical double layer and of diffusion into the interpretation of the oscillograms did yield one method which appeared to be theoretically sound but it was considered to be unduly burdensome. Because of the non-linearity of the electrode as a circuit element, graphical and Fourier analysis of the oscillographic record would be required to interpret it fully.

By making certain simplifying assumptions, including the assumption that the impedance of the double layer is essentially infinite, we developed from first principles an equation which seems to be at least roughly valid. It is:

$$\exp \frac{e_0 n F}{RT} = \frac{n F A C_{ox}}{a} \sqrt{2 \omega D_{Red}} \exp (E^0 n F / RT) - \frac{\sqrt{D_{Red}}}{\sqrt{D_{ox}}} \exp (E^0 n F / RT)$$

Where  $A$  is electrode area,  $a$  is the alternating current amplitude,  $D$  is a diffusion coefficient,  $\omega$  is the angular frequency of the alternating current,  $C_{ox}$  is the bulk concentration of the oxidant ( $C_{Red}$  being zero) and  $E^0$  is the Nernst standard potential. This equation predicts the experimentally observed shape of the curve obtained by plotting  $e_0$  against current density. When the quantity  $\exp (e_0 n F / RT)$  is plotted against  $1/a$ , the points are rather widely scattered but a straight line can be reasonably drawn through them. The frequency-dependency of  $e_0$  called for by this equation is also roughly satisfied. On the whole, the results indicate that the equation is at least roughly valid but that our data are not accurate enough for an adequate test. We propose to pursue



this line of approach further, improving our experimental accuracy and our working equation if we can.

The difficulty in interpreting the cyclograms with certainty and also the experimental difficulty of hysteresis has led us to consider an alternative alternating current electrolysis method, namely the incremental (or differential) method utilizing a small A.C. superimposed on D.C. and employing an impedance bridge technique. This is the method used by Grahame in his measurements of double layer capacities and the one assumed in his theory of the faradaic admittance.

This method appears to offer to our case several advantageous points. Hysteresis, which interfered with reproducibility when using high alternating current densities, should be considerably reduced by the incremental technique since the AC excitation must be kept negligibly small; the direct current also is generally quite small. Kinetic constants for the electrode reaction are obtainable by Grahame's theory of the Faradaic admittance directly from the bridge data and the standard potential can be calculated from the ratio of these constants. The disadvantage of this method is that it is only applicable when the electron transfer step is sufficiently slow so that the velocity constants can be measured by this method with suitable accuracy.

We have begun construction of an exceedingly well shielded bridge and suitably sensitive detector to test this method for our purposes. Since currents are in the microampere range, the shielding problem is difficult.